

Quantum mechanics without operational equivalence

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Ontological models for quantum mechanics can be different with respect to causal structure, contextuality, fine-tuning, etc. when operators are realized by different measurements

- I. Operational theories and ontological models
- II. The Bridgmannian perspective
- III. Trivialization: removing simultaneous measurements

- **Operational theory:**

- $\{P, P' \dots\}$: preparations
- $\{M, M' \dots\}$: measurements
- $\{X, X' \dots\}$: outcomes

$$\{p(X|M \wedge P) : \text{for all } M, P\}$$

- **Ontological (hidden variable) model:**

- $\{\Lambda, \Lambda' \dots\}$: ontic (hidden) states

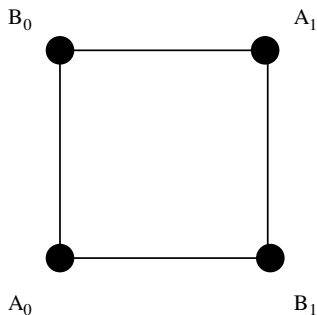
$$\{p(\Lambda|P) : \text{for all } P\}$$

$$\{p(X|M \wedge \Lambda) : \text{for all } M, \Lambda\}$$

$$p(X|M \wedge P) = \sum_{\Lambda} p(X|M \wedge \Lambda) p(\Lambda|P)$$

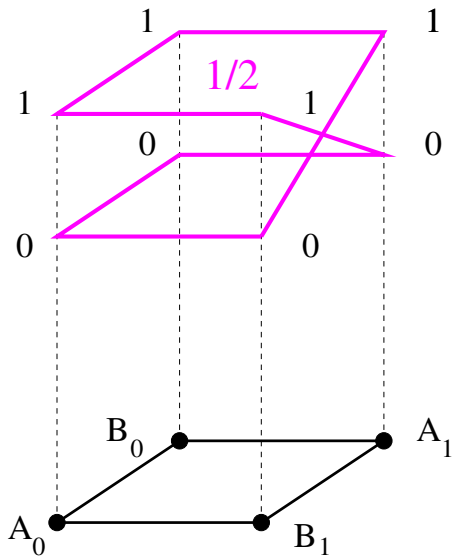
- **Measurements:**

$$\left\{ \underbrace{A_0, A_1, B_0, B_1}_{\text{Basic}}, \underbrace{A_0 \wedge B_0, A_0 \wedge B_1, A_1 \wedge B_0, A_1 \wedge B_1}_{\text{Simultaneous}} \right\}$$

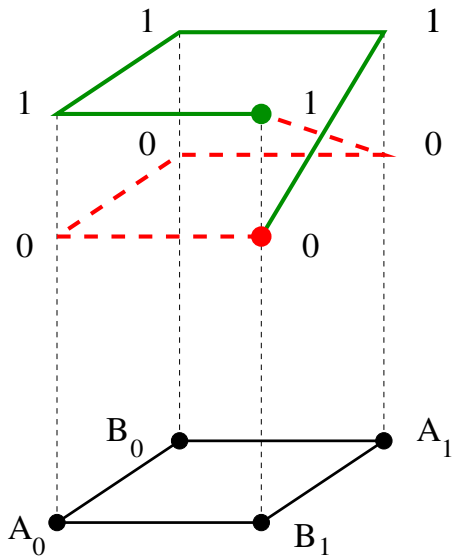


- **Outcomes:** $X, Y = 0, 1$

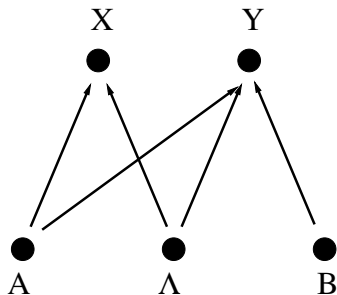
An operational theory: the PR box



An ontological model: the PR box

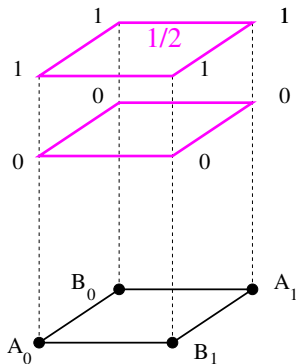


Causal structure of the PR box

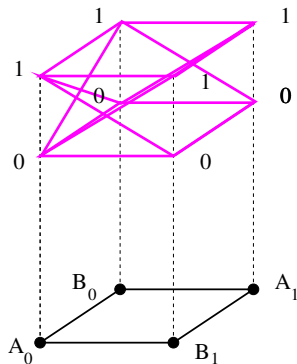


- CHSH inequality
- Contextuality
- Fine-tuning

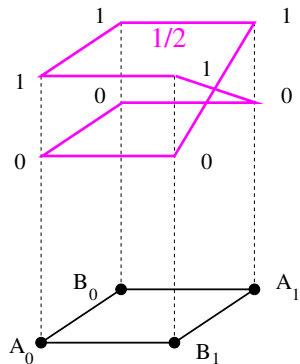
Three operational theories



Classical

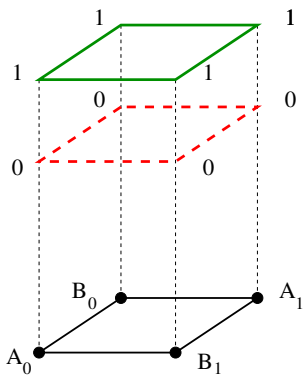


EPR-Bell

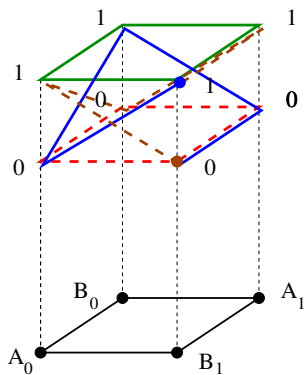


PR-box

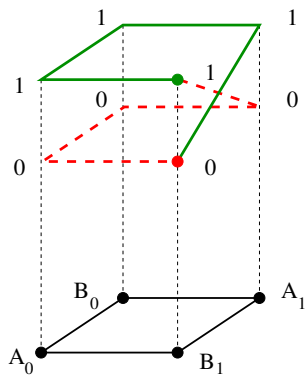
Three ontological models



Classical

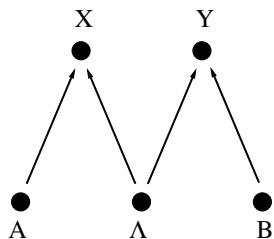


EPR-Bell

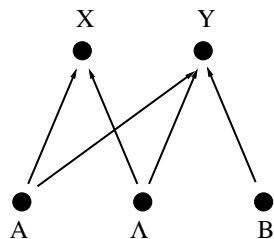


PR-box

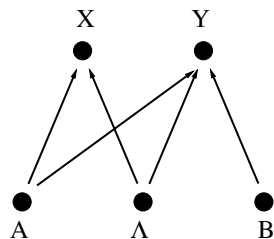
Causal structure of the models



Classical



EPR-Bell



PR box

- CHSH inequality
- Contextuality
- Fine-tuning

- **Representation:**

$$p(X|M \wedge P) = \text{Tr}(\boldsymbol{\rho} \mathbf{P})$$

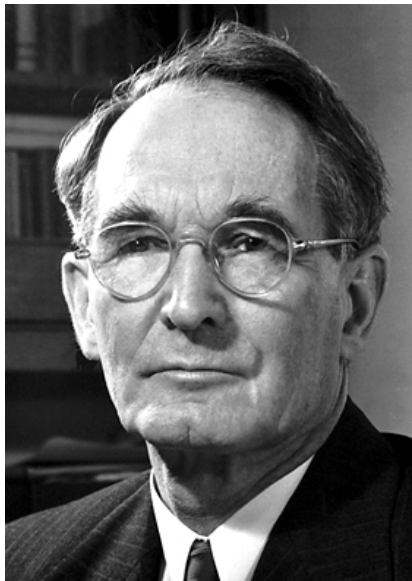
- $\{P, P' \dots\} \longrightarrow \{\boldsymbol{\rho}, \boldsymbol{\rho}' \dots\}$: density operators
- $\{M, M' \dots\} \longrightarrow \{\mathbf{O}, \mathbf{O}' \dots\}$: self-adjoint operators
- $\{X, X' \dots\} \longrightarrow \{\mathbf{P}, \mathbf{P}' \dots\}$: spectral projections

Operational equivalence

- $M_1 \sim M_2$ if $p(X|M_1 \wedge P) = p(X|M_2 \wedge P)$ for all P
- Then M_1 and M_2 are represented by the same operator:

$$p(X|M_1 \wedge P) = p(X|M_2 \wedge P) = \text{Tr}(\boldsymbol{\rho} \mathbf{P})$$

Percy Bridgman (1882-1961)



Bridgman's question

When do two measurements measure the same observable?

Bridgmannian

Standard

Operator:

\mathcal{O}

\mathcal{O}



Observable:

\mathcal{O}_1

\mathcal{O}_2

\mathcal{O}



Measurement:

M_1

\sim

M_2

M_1

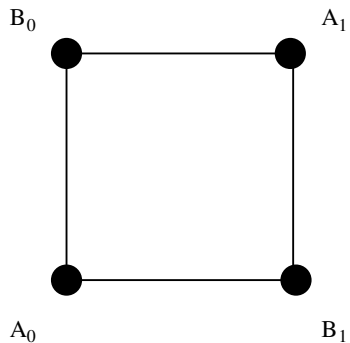
\sim

M_2

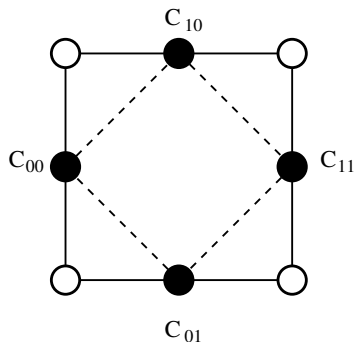
- Keep the operators, change the measurements
- Replace measurements by operationally equivalent measurements
- Remove simultaneous measurements
- Show that the new operational theory is different with respect to the causal structure, contextuality, fine-tuning, etc.

Replace the simultaneous measurements by operationally equivalent new basic measurements

Trivialization



Trivialization



- The old basic basic measurements will be operationally equivalent to certain marginalizations of the new basic measurements

Trivialization

Non-trivial theory \longrightarrow **Trivial theory**

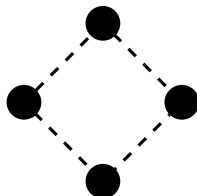
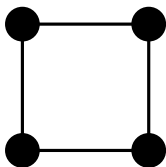
$\{A_0, A_1, B_0, B_1, A_0 \wedge B_0$ \longrightarrow $\{C_{00}, C_{01}, C_{10}, C_{11}\}$
 $A_0 \wedge B_1, A_1 \wedge B_0, A_1 \wedge B_1\}$

$$C_{00}^{(1)} \sim C_{01}^{(1)} \sim A_0$$

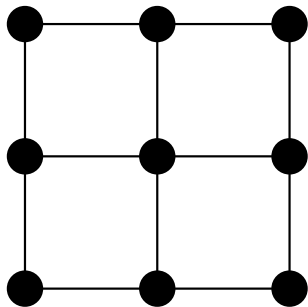
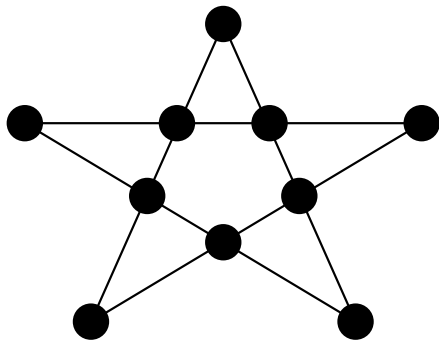
$$C_{10}^{(1)} \sim C_{11}^{(1)} \sim A_1$$

$$C_{00}^{(2)} \sim C_{10}^{(2)} \sim B_0$$

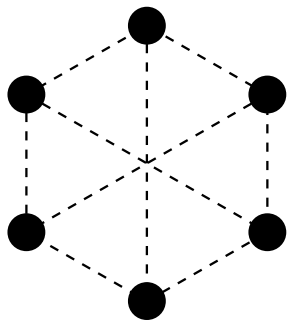
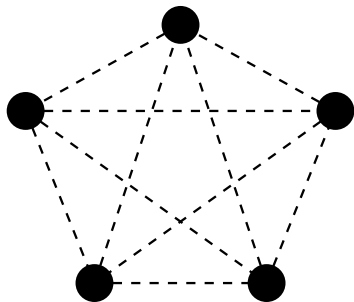
$$C_{01}^{(2)} \sim C_{11}^{(2)} \sim B_1$$



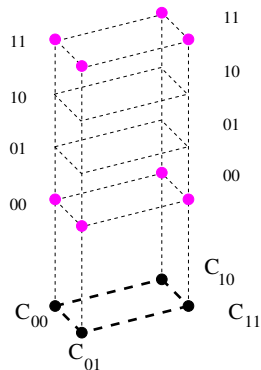
The GHZ and the Peres-Mermin graph



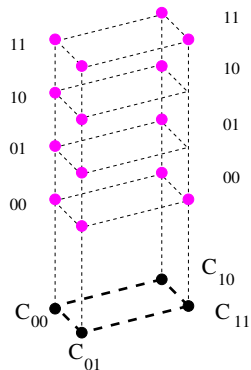
The GHZ and the Peres-Mermin line graph



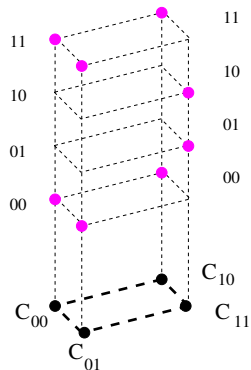
Three operational theories



Classical

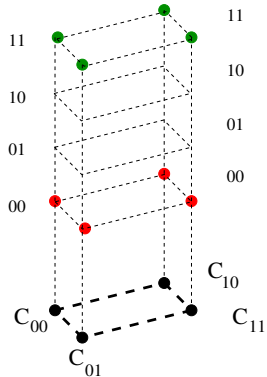


EPR-Bell

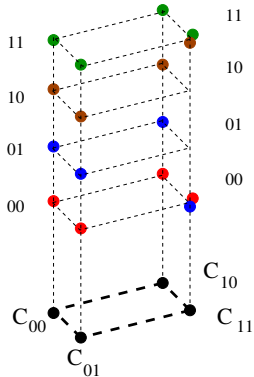


PR-box

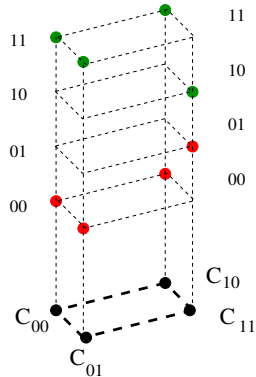
Three ontological models



Classical

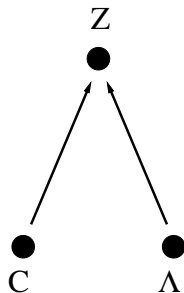


EPR-Bell



PR-box

Causal structure of the models



- CHSH inequality
- Contextuality
- Fine-tuning

$$p(X, Y|A, B, P_{\text{EPR}}) = \langle \Psi_s | (\mathbf{X}^A \otimes \mathbf{Y}^B) \Psi_s \rangle$$

Operator:

$$\mathbf{A}_0 \otimes \mathbf{B}_0$$



Measurement:

$$A_0 \wedge B_0$$

\sim

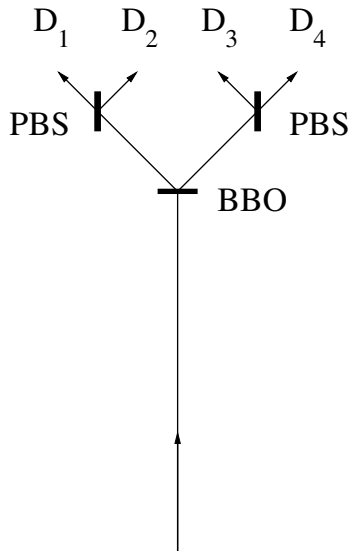
$$C_{00}$$

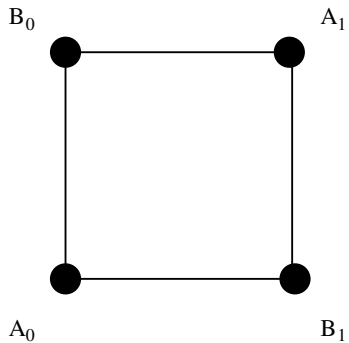
$A_0 \wedge B_0$: Measure the linear polarization of the left
photon along a transverse axis a_0

and

measure the linear polarization of the right
photon along a transverse axis b_0

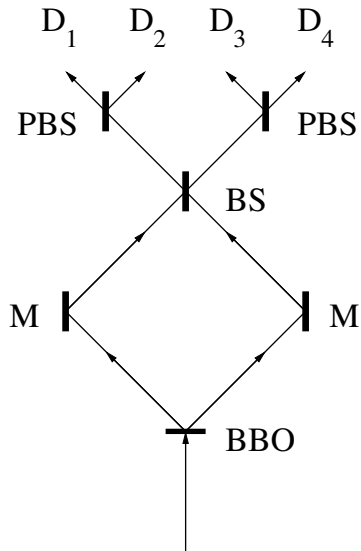
Quantum mechanics

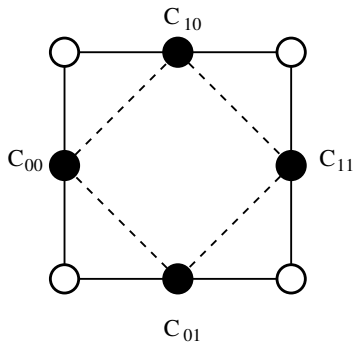




C_{00} : Perform a *global* polarization measurement on the photon pair with four outcomes corresponding to the four eigenvectors of $\mathbf{A}_0 \otimes \mathbf{B}_0$ in $H_2 \otimes H_2$

Quantum mechanics





The realization of an operator in quantum mechanics by different measurements can give rise to different ontological models with respect to contextuality, causal structure, fine-tuning, etc.

- Gábor Hofer-Szabó, "Quantum mechanics without operational equivalence" (in preparation).

Trivialization of a simple theory

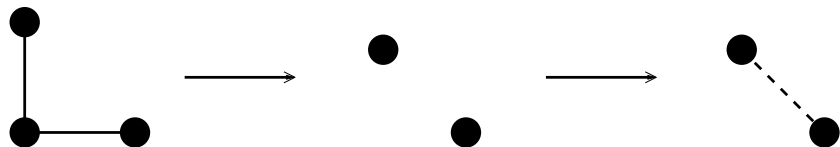
Non-trivial theory \longrightarrow **Trivial theory**

$\{M_1, M_2, M_3, M_1 \wedge M_2, M_1 \wedge M_3\}$ \longrightarrow $\{M_{12}, M_{13}\}$

$M_{12} \sim M_1 \wedge M_2$

$M_{13} \sim M_1 \wedge M_3$

$M_{12}^{(1)} \sim M_{13}^{(1)} \sim M_1$



Noncontextuality

An ontological model for QM is **noncontextual** if

- every ontic state determines the probability distribution of outcomes of every measurement independently of what other measurements are simultaneously performed
(simultaneous noncontextuality)
- any two measurements which are represented by the same self-adjoint operator have the same probability distribution of outcomes in every ontic state
(measurement noncontextuality)

- **Simultaneous noncontextuality:**

$$p(X|M \wedge \Lambda) = p(X|M \wedge M' \wedge \Lambda) \quad \text{for all } \Lambda$$

- **Measurement noncontextuality:**

$$\text{If } p(X|M \wedge P) = p(X'|M' \wedge P) \quad \text{for all } P$$

$$\text{then } p(X|M \wedge \Lambda) = p(X'|M' \wedge \Lambda) \quad \text{for all } \Lambda$$